

Review the simple (haha!) case of number theory over \mathbb{Z} :

Continuing discussion of analytical properties of $\zeta(s)$ relevant to *Riemann's Explicit Formula* (von Mangoldt's reformulation):

$$\sum_{p^m < X} \log p = X - (b+1) - \lim_{T \rightarrow \infty} \sum_{|\operatorname{Im}(\rho)| < T} \frac{X^\rho}{\rho} + \sum_{n \geq 1} \frac{X^{-2n}}{2n}$$

We are in the course of proving that the *completed* zeta function

$$\xi(s) = \pi^{-\frac{s}{2}} \Gamma\left(\frac{s}{2}\right) \zeta(s)$$

has an analytic continuation to $s \in \mathbb{C}$, except for simple poles at $s = 0, 1$, and has the *functional equation*

$$\xi(1-s) = \xi(s)$$

... and (anticipating the Riemann-Hadamard product issues) $s(s-1)\xi(s)$ is entire and **bounded in vertical strips**.

We need the simplest *theta function*

$$\theta(z) = \sum_{n \in \mathbb{Z}} e^{\pi i n^2 z} \quad (\text{with } z \in \mathfrak{H})$$

By Riemann's time, Jacobi's functional equation of $\theta(z)$ was well-known, as the simplest example of a larger thing:

$$\theta(z) = \frac{1}{\sqrt{-iz}} \cdot \theta(-1/z)$$

(Proof below.) The modified version

$$\frac{\theta(iy) - 1}{2} = \sum_{n=1}^{\infty} e^{-\pi n^2 y}$$

gets used just below.

The connection to $\zeta(s)$ is the *integral presentation*:

Claim: For $\operatorname{Re}(s) > 1$

$$\pi^{-s/2} \Gamma\left(\frac{s}{2}\right) \cdot \zeta(s) = \int_0^\infty \frac{\theta(iy) - 1}{2} \cdot y^{s/2} \cdot \frac{dy}{y}$$

Meaning? An integral against t^s with dt/t , a *Mellin transform*, is just a *Fourier transform* in different coordinates.

Starting from the integral, for $\operatorname{Re}(s) > 1$, compute directly

$$\begin{aligned} \int_0^\infty \frac{\theta(iy) - 1}{2} y^{s/2} \frac{dy}{y} &= \int_0^\infty \sum_{n \geq 1} e^{-\pi n^2 y} y^{s/2} \frac{dy}{y} \\ &= \sum_{n \geq 1} \int_0^\infty e^{-\pi n^2 y} y^{s/2} \frac{dy}{y} = \pi^{-s/2} \sum_{n \geq 1} \frac{1}{n^{2s}} \int_0^\infty e^{-y} y^{s/2} \frac{dy}{y} \end{aligned}$$

by replacing y by $y/(\pi n^2)$, and interchanging sum and integral, giving

$$= \pi^{-s/2} \Gamma\left(\frac{s}{2}\right) \cdot \sum_{n \geq 1} \frac{1}{n^s} = \zeta(s) \quad (\text{for } \operatorname{Re}(s) > 1)$$

$\frac{\theta(iy) - 1}{2} = \sum_{n=1}^{\infty} e^{-\pi n^2 y}$ is of rapid decay as $y \rightarrow +\infty$:

$$\begin{aligned} \frac{\theta(iy) - 1}{2} &= \sum_{n \geq 1} e^{-\pi n^2 y} \leq e^{-\pi y/2} \sum_{n \geq 1} e^{-\pi n^2/2} \\ &= \text{const} \cdot e^{-\pi y/2} \quad (\text{for } y \geq 1) \end{aligned}$$

Thus, the integral from 1 (not 0) to $+\infty$ is nicely convergent for *all* values of s , and

$$\int_1^{\infty} \frac{\theta(iy) - 1}{2} y^{s/2} \frac{dy}{y} = \text{entire in } s$$

The trick (known before Riemann) is to use Jacobi's functional equation for $\theta(z)$ to convert the part of the integral from 0 to 1 into a similar integral from 1 to $+\infty$.

It is not obvious that $\theta(iy)$ has any property that would ensure this. However, in the early 19th century theta functions were intensely studied.

Again, the functional equation of θ , proven below, is

$$\theta(z) = \frac{1}{\sqrt{-iz}} \cdot \theta(-1/z)$$

Book-keeping:

$$\frac{\theta(-1/iy) - 1}{2} = y^{1/2} \frac{\theta(iy) - 1}{2} + \frac{y^{1/2}}{2} - \frac{1}{2}$$

Then

$$\begin{aligned} \int_0^1 \frac{\theta(iy) - 1}{2} y^{s/2} \frac{dy}{y} &= \int_1^\infty \frac{\theta(-1/iy) - 1}{2} y^{-s/2} \frac{dy}{y} \\ &= \int_1^\infty \left(y^{1/2} \frac{\theta(iy) - 1}{2} + \frac{y^{1/2}}{2} - \frac{1}{2} \right) y^{-s/2} \frac{dy}{y} \\ &= \int_1^\infty \frac{\theta(iy) - 1}{2} y^{-s/2} \frac{dy}{y} + \int_1^\infty \left(\frac{y^{(1-s)/2}}{2} - \frac{y^{-s/2}}{2} \right) \frac{dy}{y} \\ &= \int_1^\infty \frac{\theta(iy) - 1}{2} y^{-s/2} \frac{dy}{y} + \frac{1}{s-1} - \frac{1}{s} \\ &= (\text{entire}) + \frac{1}{s-1} - \frac{1}{s} \end{aligned}$$

The elementary expressions $1/(s - 1)$ and $1/s$ certainly have meromorphic continuations to \mathbb{C} , with explicit poles. Thus, together with the first integral from 1 to ∞ , we have

$$\begin{aligned} & \pi^{-s/2} \Gamma\left(\frac{s}{2}\right) \zeta(s) \\ &= \int_1^\infty \frac{\theta(iy) - 1}{2} (y^{s/2} + y^{(1-s)/2}) \frac{dy}{y} + \frac{1}{s-1} - \frac{1}{s} \\ &= (\text{entire}) + \frac{1}{s-1} - \frac{1}{s} \end{aligned}$$

The right-hand side is visibly symmetrical under $s \rightarrow 1 - s$, which gives the functional equation. ///

Comments: Attempting to avoid the gamma factor $\pi^{-s/2} \Gamma\left(\frac{s}{2}\right)$ leads to an unsymmetrical and unenlightening form.

The fact that $\Gamma(s/2)$ has no zeros assures that it masks no poles of $\zeta(s)$. Non-vanishing of $\Gamma(s)$ follows from the identity

$$\Gamma(s) \cdot \Gamma(1 - s) = \frac{\pi}{\sin \pi s}$$

Claim: Jacobi's functional equation for $\theta(z)$

$$\theta(-1/iy) = \sqrt{y} \cdot \theta(iy)$$

Proof: This symmetry itself follows from a more fundamental fact, the **Poisson summation formula**

$$\sum_{n \in \mathbb{Z}} f(n) = \sum_{n \in \mathbb{Z}} \widehat{f}(n) \quad (\widehat{f} \text{ is Fourier transform})$$

$$\text{Fourier transform of } f = \widehat{f}(\xi) = \int_{\mathbb{R}} f(x) e^{-2\pi i x \xi} dx$$

The Poisson summation formula is applied to

$$f(x) = \varphi(\sqrt{y} \cdot x) \quad \text{with} \quad \varphi(x) = e^{-\pi x^2}$$

The Gaussian $\varphi(x) = e^{-\pi x^2}$ has the useful property that it is its own Fourier transform.

Prove that the Gaussian is its own Fourier transform by completing the square and a contour integration shift:

$$\begin{aligned}\widehat{\varphi}(\xi) &= \int_{\mathbb{R}} e^{-\pi x^2} e^{-2\pi i x \xi} dx = \int_{\mathbb{R}} e^{-\pi(x+i\xi)^2 - \pi\xi^2} dx \\ &= e^{-\pi\xi^2} \int_{\mathbb{R}} e^{-\pi(x+i\xi)^2} dx\end{aligned}$$

By moving the contour of integration, the latter integral is

$$\int_{\mathbb{R}} e^{-\pi(x+i\xi)^2} dx = \int_{\mathbb{R}+i\xi} e^{-\pi x^2} dx = \int_{\mathbb{R}} e^{-\pi x^2} dx$$

Thus, the integral is independent of ξ . In fact, the constant is 1. By a straightforward change of variables, Fourier transform behaves well with respect to *dilations*:

$$\begin{aligned}\widehat{f}(\xi) &= \int_{\mathbb{R}} \varphi(\sqrt{y} x) e^{-2\pi i x \xi} dx = \frac{1}{\sqrt{y}} \int_{\mathbb{R}} \varphi(x) e^{-2\pi i x \xi / \sqrt{y}} dx \\ &= \frac{1}{\sqrt{y}} \widehat{\varphi}(\xi / \sqrt{y}) = \frac{1}{\sqrt{y}} e^{-\pi \xi^2 / y} \quad (\text{replacing } x \text{ by } x / \sqrt{y})\end{aligned}$$

Applying Poisson summation to $f(x) = e^{-\pi x^2 y}$,

$$\sum_{n \in \mathbb{Z}} e^{-\pi n^2 y} = \frac{1}{\sqrt{y}} \sum_{n \in \mathbb{Z}} e^{-\pi n^2 / y}$$

This gives

$$\theta(iy) = \frac{1}{\sqrt{y}} \theta(-1/iy)$$

Remark For $z \in \mathfrak{H}$, also $-1/z \in \mathfrak{H}$, and the series for $\theta(z)$ and $\theta(-1/z)$ are nicely convergent. The identity proven for θ is $\theta(-1/z) = \sqrt{-iz} \theta(z)$ on the imaginary axis. The *Identity Principle* from complex analysis implies that the same equality holds for all $z \in \mathfrak{H}$.

Heuristic for **Poisson summation** $\sum_{n \in \mathbb{Z}} f(n) = \sum_{n \in \mathbb{Z}} \widehat{f}(n)$

The periodicized version of a function f on \mathbb{R} is

$$F(x) = \sum_{n \in \mathbb{Z}} f(x + n)$$

A periodic function should be (!) represented by its *Fourier series*:

$$F(x) = \sum_{\ell \in \mathbb{Z}} e^{2\pi i \ell x} \int_0^1 F(x) e^{-2\pi i \ell x} dx$$

Fourier *coefficients* of F expand to be the Fourier *transform* of f :

$$\begin{aligned} \int_0^1 F(x) e^{-2\pi i \ell x} dx &= \int_0^1 \sum_{n \in \mathbb{Z}} f(x + n) e^{-2\pi i \ell x} dx \\ &= \sum_{n \in \mathbb{Z}} \int_n^{n+1} f(x) e^{-2\pi i \ell x} dx = \int_{\mathbb{R}} f(x) e^{-2\pi i \ell x} dx = \widehat{f}(\ell) \end{aligned}$$

Evaluating at 0, we should have

$$\sum_{n \in \mathbb{Z}} f(n) = F(0) = \sum_{\ell \in \mathbb{Z}} \widehat{f}(\ell)$$

What would it take to legitimize this?

Certainly f must be of sufficient decay so that the integral for its Fourier transform is convergent. and so that summing its translates by \mathbb{Z} is convergent.

We'd want f to be continuous, probably differentiable, so that we can talk about pointwise values of F

... and to make plausible the hope that the Fourier series of F converges to F pointwise.

For f and several derivatives rapidly decreasing, the Fourier transform \widehat{f} will be of sufficient decay so that its sum over \mathbb{Z} does converge.

A simple sufficient hypothesis for convergence is that f be in the *Schwartz space* of infinitely-differentiable functions all of whose derivatives are of *rapid decay*, that is,

Schwartz space = $\{ \text{smooth } f : \sup_x (1 + x^2)^\ell |f^{(i)}(x)| < \infty \text{ for all } i, \ell \}$

Representability of a periodic function by its Fourier series is a serious question, with several possible senses. We want *pointwise convergence*. A special, self-contained argument gives a good-enough result for immediate purposes.

Consider (\mathbb{Z} -)periodic functions on \mathbb{R} , that is, complex-valued functions f on \mathbb{R} such that $f(x + n) = f(x)$ for all $x \in \mathbb{R}$, $n \in \mathbb{Z}$. For periodic f sufficiently nice so that integrals

$$\widehat{f}(n) = \int_0^1 f(x) e^{-2\pi i n x} dx \quad (n^{\text{th}} \text{ Fourier coefficient of } f)$$

make sense, the **Fourier expansion** of f is

$$f \sim \sum_{n \in \mathbb{Z}} \widehat{f}(n) e^{2\pi i n x}$$

We want

$$f(x_o) = \sum_{n \in \mathbb{Z}} \widehat{f}(n) e^{2\pi i n x_o}$$

Consider periodic *piecewise- C^0* functions which are left-continuous and right-continuous at any discontinuities.

Theorem: For periodic piecewise- C^0 function f , left-continuous and right-continuous at discontinuities, for points x_o at which f is C^0 and *left-differentiable* and *right-differentiable*, the Fourier series of f evaluated at x_o converges to $f(x)$:

$$f(x_o) = \sum_{n \in \mathbb{Z}} \hat{f}(n) e^{2\pi i n x_o}$$

That is, for such functions, at such points, the Fourier series *represents* the function *pointwise*.

A notable missing conclusion is *uniform* pointwise convergence. For more serious applications, pointwise convergence not known to be uniform is often useless.

Proof: Can reduce to $x_0 = 0$ and $f(0) = 0$. Representability of $f(0)$ by the Fourier series is the assertion that

$$\begin{aligned} 0 = f(0) &= \lim_{M, N \rightarrow +\infty} \sum_{-M \leq n < N} \widehat{f}(n) e^{2\pi i n \cdot 0} \\ &= \lim_{M, N \rightarrow +\infty} \sum_{-M \leq n < N} \widehat{f}(n) \end{aligned}$$

Substituting the defining integral for the Fourier coefficients:

$$\begin{aligned} \sum_{-M \leq n < N} \widehat{f}(n) &= \sum_{-M \leq n < N} \int_0^1 f(u) e^{-2\pi i n u} du \\ &= \int_0^1 \sum_{-M \leq n < N} f(u) e^{-2\pi i n u} du = \int_0^1 f(u) \cdot \frac{e^{2\pi i M u} - e^{-2\pi i N u}}{1 - e^{-2\pi i u}} du \end{aligned}$$

We will show that

$$\lim_{\ell \rightarrow \pm\infty} \int_0^1 \frac{f(u) \cdot e^{-2\pi i \ell u}}{1 - e^{-2\pi i u}} du = 0$$

Since $f(0) = 0$, the function

$$g(x) = \frac{f(x)}{1 - e^{-2\pi ix}}$$

is piecewise- C^∞ , and left-continuous and right-continuous at discontinuities. The only issue is at integers, and by the periodicity it suffices to prove continuity at 0.

$$\frac{f(x)}{1 - e^{-2\pi ix}} = \frac{f(x)}{x} \cdot \frac{x}{1 - e^{-2\pi ix}}$$

The two-sided limit

$$\lim_{x \rightarrow 0} \frac{x}{1 - e^{-2\pi ix}} = \left. \frac{d}{dx} \right|_{x=0} \frac{x}{1 - e^{-2\pi ix}}$$

exists, by differentiability. Similarly, we have left and right limits

$$\lim_{x \rightarrow 0^-} \frac{f(x)}{x} \quad \text{and} \quad \lim_{x \rightarrow 0^+} \frac{f(x)}{x}$$

by the one-sided differentiability of f . So both one-sided limits exist, giving the one-sided continuity of g at 0. ///

We find ourselves wanting a *Riemann-Lebesgue lemma*, that the Fourier coefficients of a periodic, piecewise- C^0 function g , with left and right limits at discontinuities, go to 0.

The essential property approximability by step functions: given $\varepsilon > 0$ there is a *step function* $s(x)$ such that

$$\int_0^1 |s(x) - g(x)| dx < \varepsilon$$

With such s ,

$$|\widehat{s}(n) - \widehat{g}(n)| \leq \int_0^1 |s(u) - g(u)| du < \varepsilon \quad (\text{for all } \varepsilon > 0)$$

It suffices that Fourier coefficients of *step functions* go to 0, an easy computation:

$$\int_a^b e^{-2\pi i \ell x} dx = \left[\frac{e^{-2\pi i \ell x}}{-2\pi i \ell} \right]_a^b = \frac{e^{-2\pi i \ell b} - e^{-2\pi i \ell a}}{-2\pi i \ell} \rightarrow 0$$

as $\ell \rightarrow \pm\infty$. Thus, the Fourier coefficients of g go to 0, so the Fourier series of f converges to $f(0)$ when f is C^1 at 0.
