

Lecture 26: Teichmüller Theory

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Math 8680 (4-25-11)

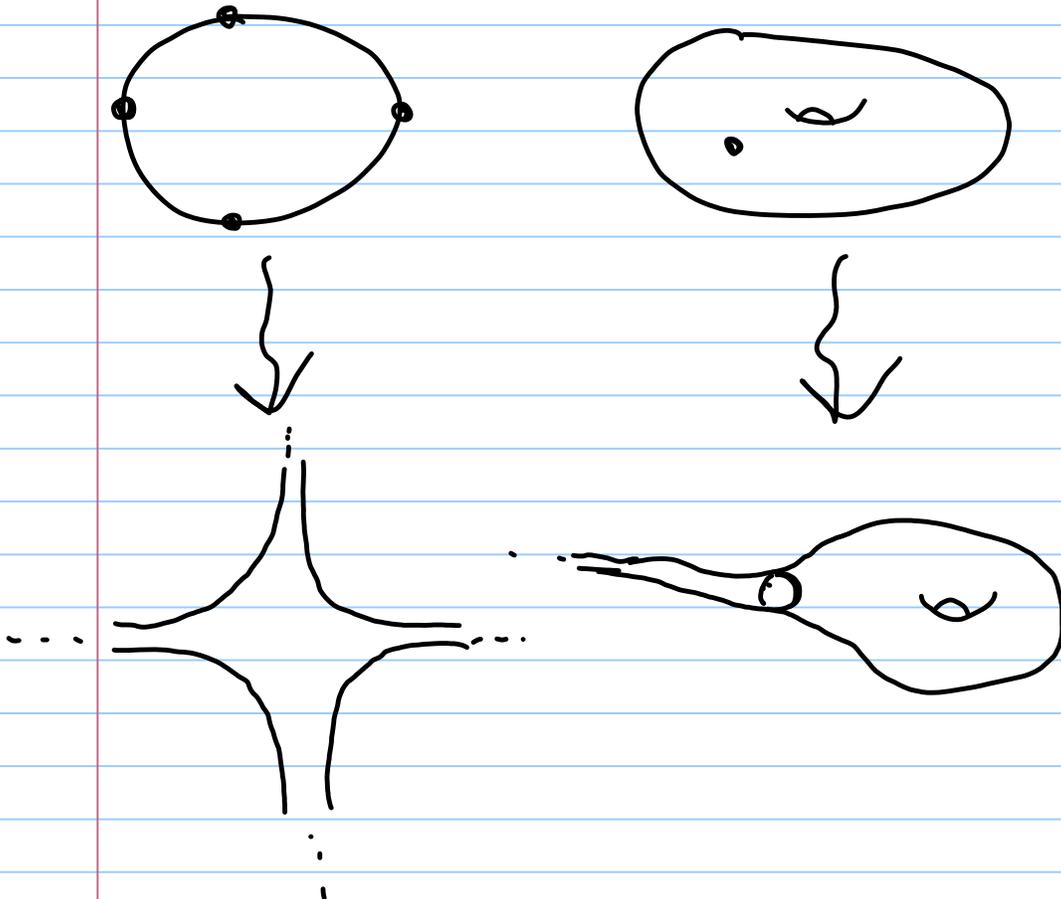
Note Title

4/25/2011

① Given a marked surface (S, M) $\mathcal{T}(S, M)$ Teichmüller is defined to be the space of metrics on (S, M) satisfying the following properties:

- hyperbolic (constant curvature -1)
- have geodesic boundary on ∂S
- have cusps at marked points M .

Examples



Fact: $\mathcal{T}(S, M)$ is a manifold of \mathbb{R} -dim $6g - 6 + 2p + 3b + c$

$\#M \cap (S - \partial S)$ \rightarrow $\#$ boundary comps. \rightarrow $\#M \cap \partial S$

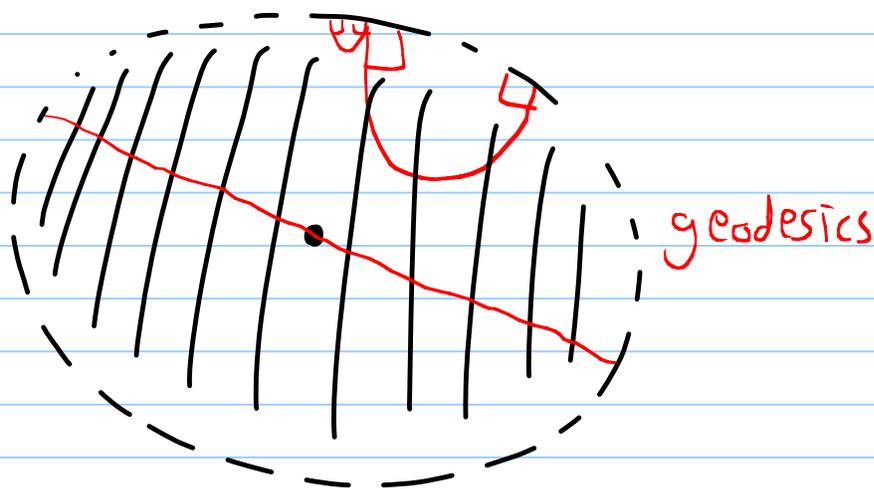
② polygon $(n+3)$ -gon
e.g. $g=0, p=0, b=1, c=n+3$

$\mathcal{J}(S, M)$ has $\dim = n$

Can view this polygon as Poincaré disk model for hyperbolic space.

Metric = ds satisfying

$$ds^2 = \frac{dx^2 + dy^2}{(1-r^2)^2} \quad \text{with} \quad r = \sqrt{x^2 + y^2}$$



Points on the boundary are infinitely far away from the center.

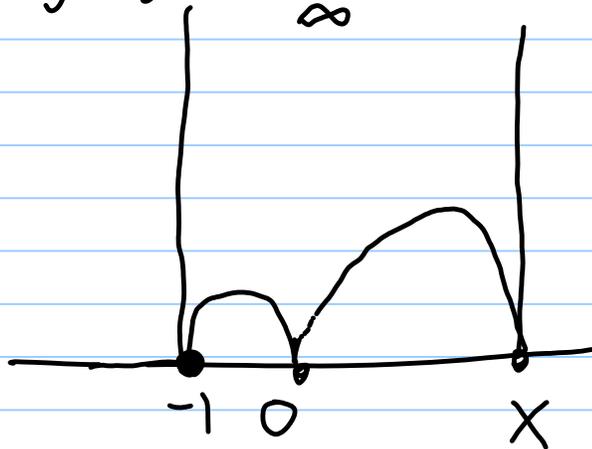
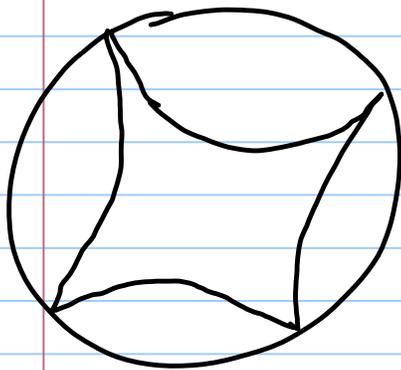
e.g. the distance from

$(1,0)$ to $(0,0)$ is

$$\int_0^1 \frac{dx}{1-x^2} = \operatorname{arctanh}(x) \Big|_0^1 \\ = \frac{1}{2} \log \left(\frac{1+x}{1-x} \right) \Big|_0^1 \rightarrow \infty.$$

③ Up to action of $PSL_2(\mathbb{R})$, we can fix 3 points on boundary and choose the other n freely.

If instead of Poincaré Disk model, we used upper half-plane, 3 points could be at $-1, 0, \infty$

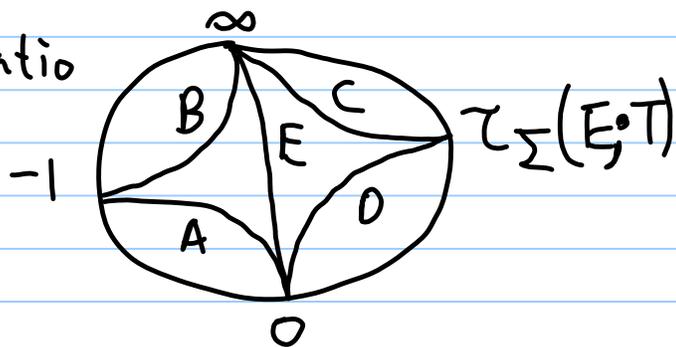


Remaining corner of a quadrilateral is free $X \in \mathbb{R}_{>0}$ but is an invariant of the quadrilateral, cross-ratio

shear coordinates.

Given a hyperbolic structure $\Sigma \in \mathcal{T}(S, M)$ and a triangulation $T = \{E_i\}_{i=1}^n$, the shear coord.

$\tau_\Sigma(E; T)$ of edge $E \in T$ is the cross-ratio

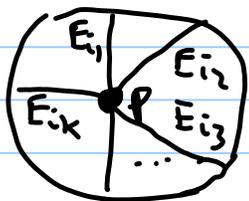


④ Theorem: The map $J(S, M) \rightarrow \mathbb{R}^n$

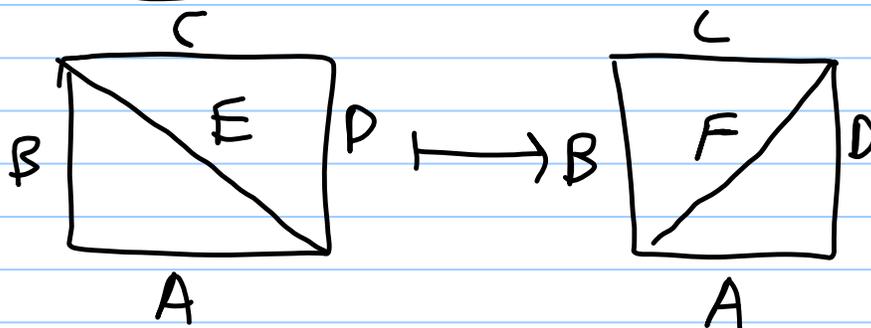
$$\Sigma \mapsto \left\{ \tau_{\Sigma}(E_{ij}; T) \right\}_{i=1}^n$$

is a homeomorphism onto the subset of \mathbb{R}^n where for each puncture p , and incident arcs $E_{i,j}, \dots, E_{i,k}$, we have

$$\prod_{j=1}^k \tau_{\Sigma}(E_{i,j}; T) = 1.$$



When we flip quads to get from $T \rightarrow T'$,



shear coordinates change in a predictable way:

$$\tau_{\Sigma}(F; T') = \tau_{\Sigma}(E; T)^{-1}$$

$$\tau_{\Sigma}(A; T') = \tau_{\Sigma}(A; T) \left(1 + \tau_{\Sigma}(E; T)^{-1} \right)^{-1}$$

$$\tau_{\Sigma}(B; T') = \tau_{\Sigma}(B; T) \left(1 + \tau_{\Sigma}(E; T) \right)$$

$$\tau_{\Sigma}(C; T') = \tau_{\Sigma}(C; T) \left(1 + \tau_{\Sigma}(E; T)^{-1} \right)^{-1}$$

$$\tau_{\Sigma}(D; T') = \tau_{\Sigma}(D; T) \left(1 + \tau_{\Sigma}(E; T) \right)$$

⑤ As we will see shortly, these mutations also appear in cluster algs!

Recall, the original definition

$(\underline{X}, \underline{Y}, B)$ $\underline{X} = \{x_{1, \dots, j} x_n\}$ initial cluster,

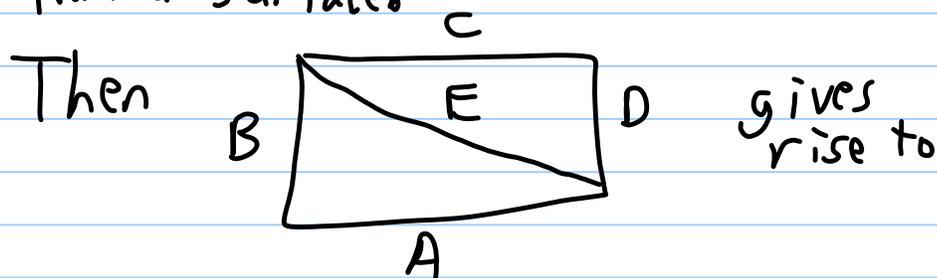
$\underline{Y} = \{y_{1, \dots, j} y_n\}$ initial coeffs, $B = n \times n$ exchange matrix

$$x'_k x_k = y_k \prod_{b_{ik} > 0} x_i^{b_{ik}} + \prod_{b_{ik} < 0} x_i^{-b_{ik}}$$

$(y_k \oplus 1)$

$$y'_j = \begin{cases} y_k^{-1} & \text{if } j=k, \\ y_j \left(\frac{y_k}{y_k \oplus 1} \right)^{b_{kj}} & \text{if } j \neq k \ \& \ b_{kj} > 0, \\ y_j (y_k \oplus 1)^{-b_{kj}} & \text{if } j \neq k \ \& \ b_{kj} \leq 0. \end{cases}$$

Let us focus on case of tropical semifield and cluster algebra from a surface.



portion of exchange matrix :

⑥

E	A	B	C	D
0	1	-1	1	-1
-1	0		*	
1		0	*	
-1	*	0		
1				0

Thus, if we mutate μ_E ,
 $X_{E'} = X_F$ and

$$Y_F = Y_{E'} = Y_E^{-1}$$

$$Y_A' = Y_A \left(\frac{Y_E}{Y_E \oplus 1} \right) = Y_A \frac{Y_E}{Y_E} \left(\frac{1}{1 \oplus 1/Y_E} \right)$$

$$Y_B' = Y_B (Y_E \oplus 1)$$

$$Y_C' = Y_C \left(\frac{Y_E}{Y_E \oplus 1} \right) = Y_C \frac{Y_E}{Y_E} \left(\frac{1}{1 \oplus 1/Y_E} \right)$$

$$Y_D' = Y_D (Y_E \oplus 1)$$

Moral: Coeff Dynamics behave like Tropical Shear coordinate dynamics.

This motivated use of laminations.

Def: An elementary lamination is like an arc except instead of going between marked points, its endpoints intersect the boundary (away from a marked point) or spiral into a puncture, or is a closed curve.

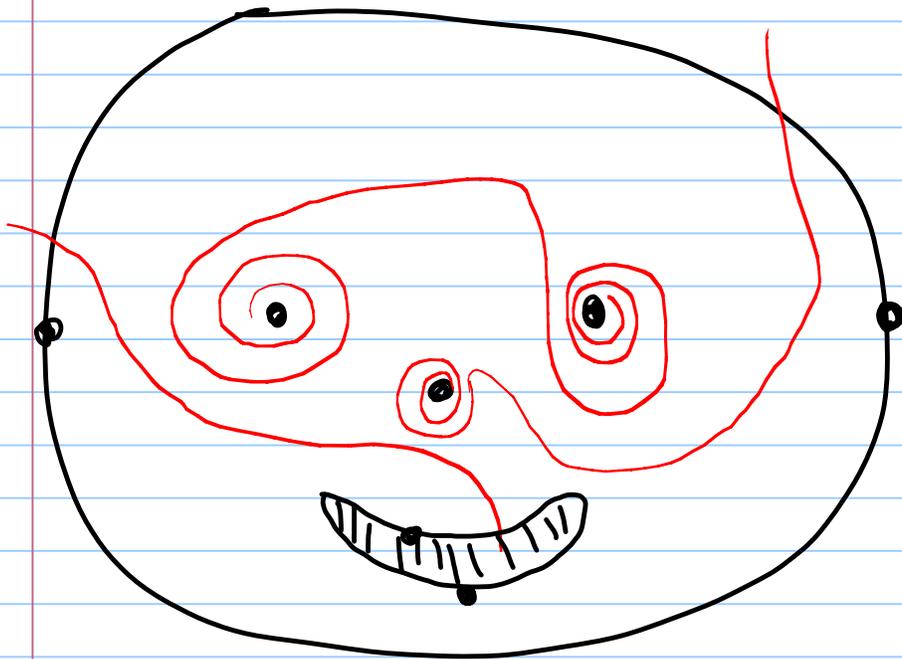
More precisely, an integral unbounded measured lamination on a marked surface (S, M) is a finite collection of pairwise non-intersecting curves, each of which has no self-intersections, modulo isotopy (relative to M) such that:

⑦ Each curve must be either :

- i) a closed curve
- ii) a curve connecting two unmarked points on ∂S
- iii) a curve with one endpoint on $\partial S - M$ and one endpoint spiraling into a puncture (clockwise OR counter-clockwise), or
- iv) a curve with both endpoints spiraling into a puncture.

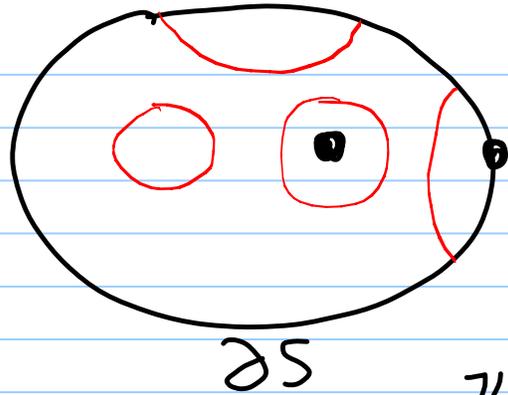
The following curves are disallowed:

- i) a closed curve bounding an unpunctured or once punctured disc,
- ii) a curve with two endpoints on ∂S which is isotopic to a boundary arc containing 0 or 1 marked point



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Disallowed

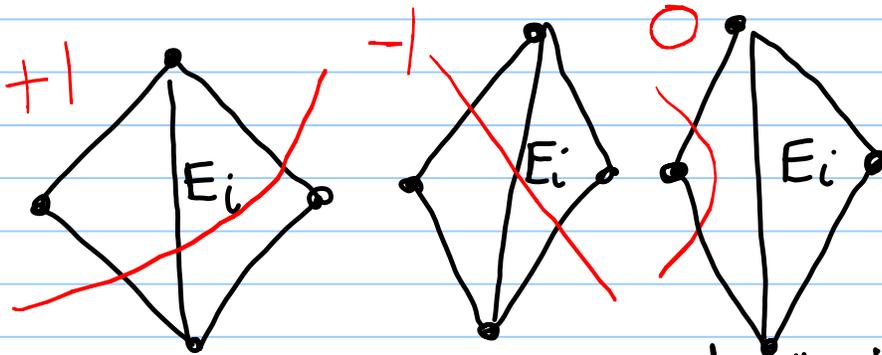


We assign a Shear coordinate E_i^L to each arc $E_i \in T$ of a triangulation w.r.t. a choice of lamination:

$$b_{E_i}(T, L) \text{ for each } E_i \in T.$$

As above, we look at quadrilateral inscribing E_i (in triangulation L)

for each curve of lamination L cutting through the quadrilateral, we calculate a contribution to the shear coordinate. Adding them all up gives contribution

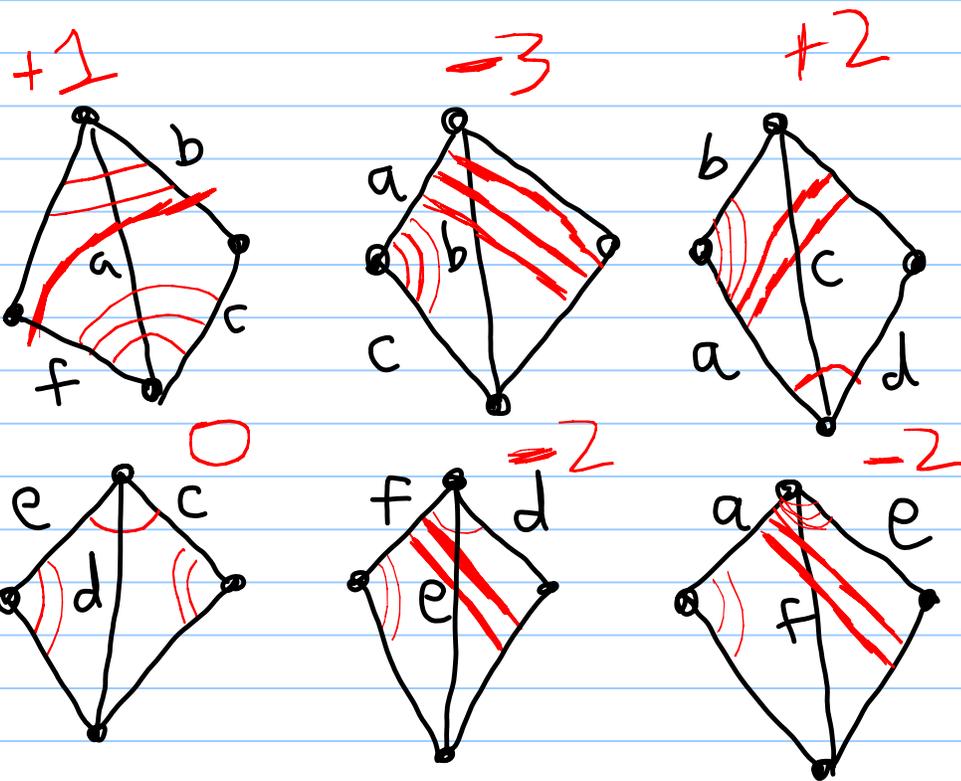
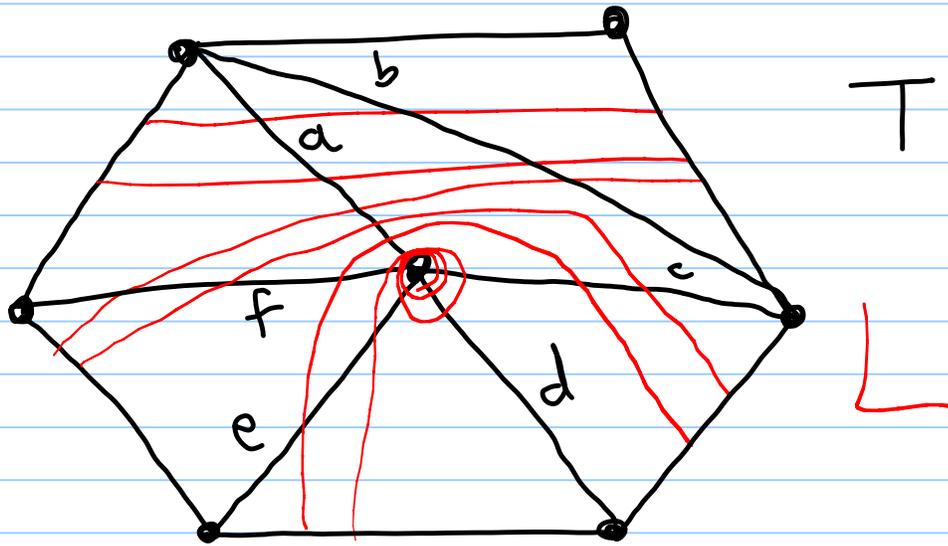


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Zilch

and all other crossings of adjacent sides give zero too.

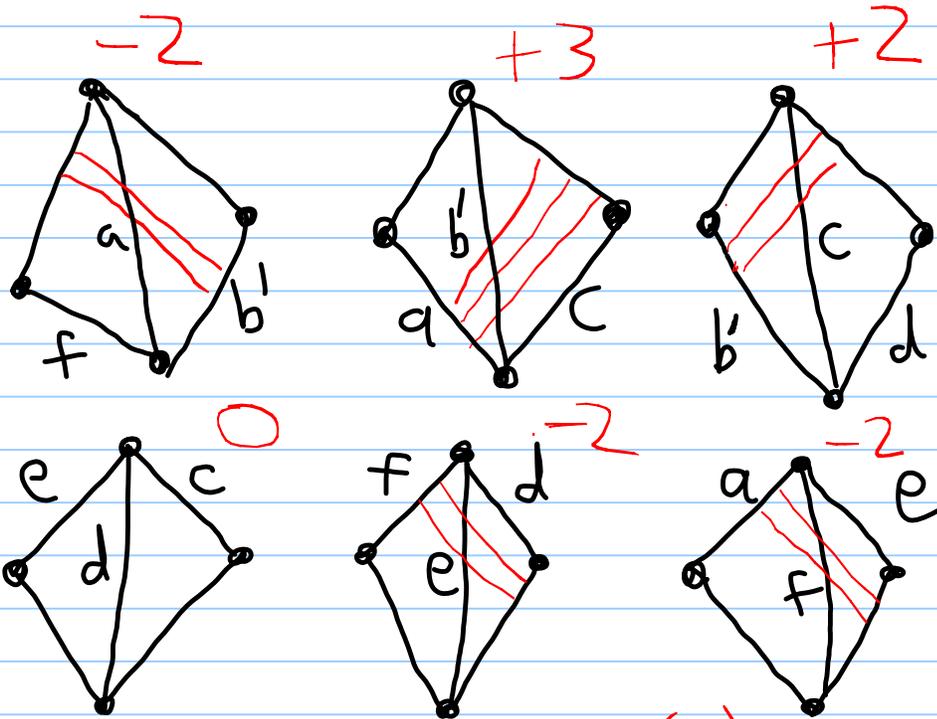
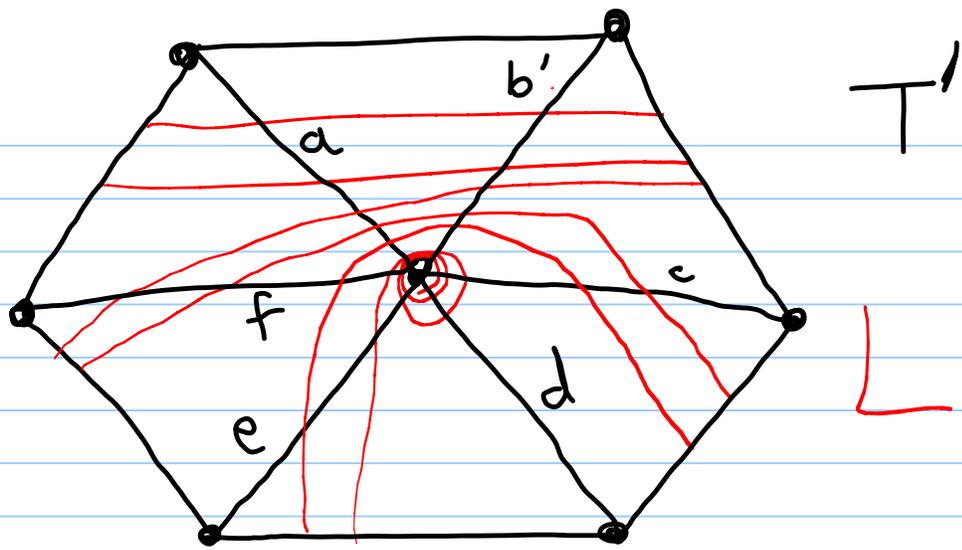
⑨ Example (Fig 3) of Fomin-Thurston)



$$b_a(T, L) = 1, b_b(T, L) = -3, \dots$$

Let us now flip $b \mapsto b'$
to get triangulation T' with
 L fixed:

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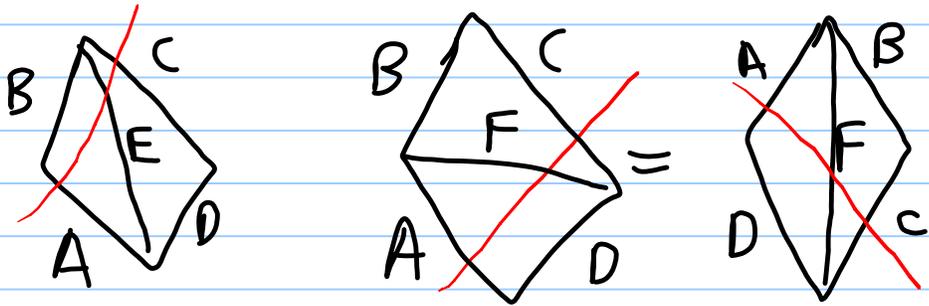
Notice: $b_{b'}(T', L) = -b_b(T, L)$

$$b_a(T', L) = b_a(T, L) - \max(-b_b(T, L), 0)$$

$$b_c(T', L) = b_c(T, L) + \max(b_b(T, L), 0)$$

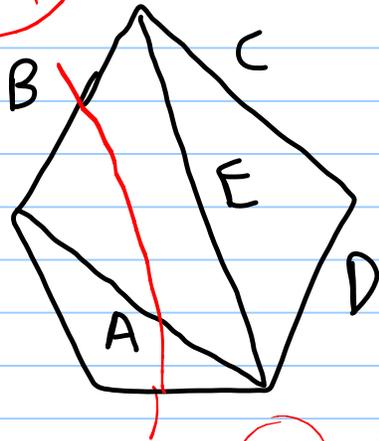
$$b_{\tau}(T', L) = b_{\tau}(T, L) \text{ o.w.}$$

(11) Not just this example, but logic works:



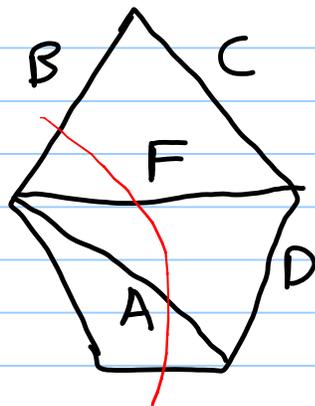
← sign reversal →

(+)



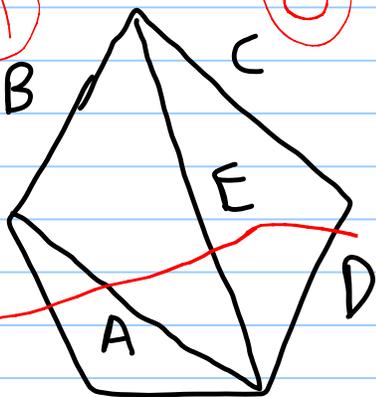
(+)

still



versus

(-)

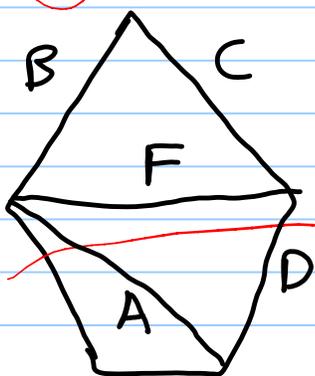


(0) w.r.t. E

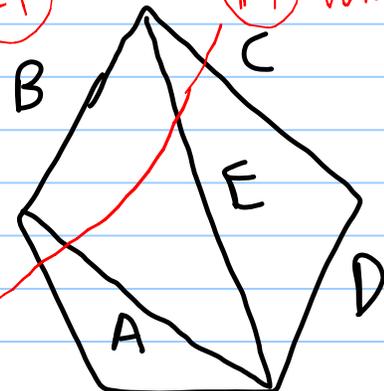
(-)

still

versus



(-)



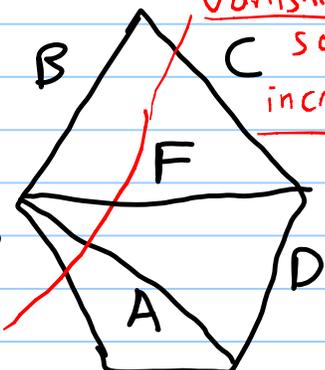
(+) w.r.t. E

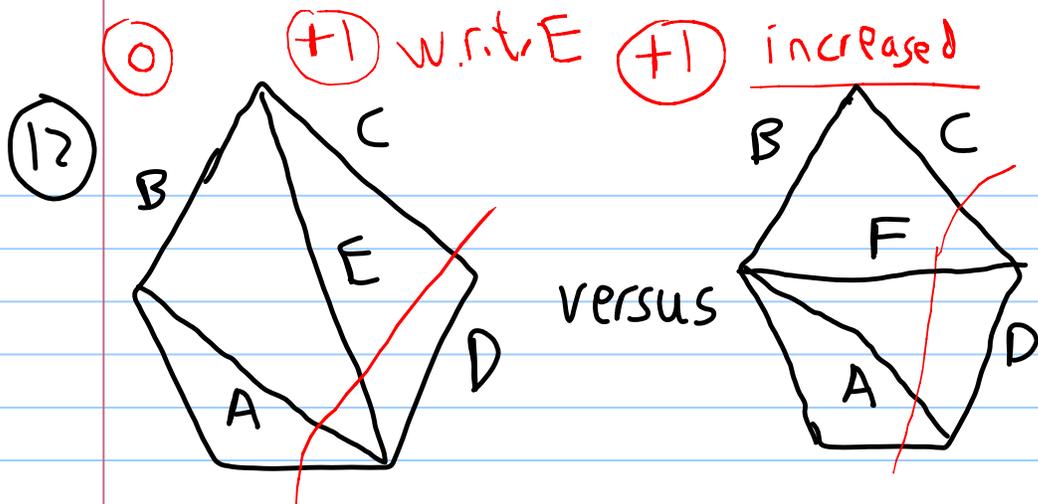
(0) contribution

vanishes

so increased

versus





$$b_A(T', L) = b_A(T, L) + \underbrace{b_E(T, L)}_{\text{if } > 0}$$

$$= b_A(T, L) + \max(b_E(T, L), 0).$$

Moral: if we let $\oplus = \max(-, -)$, coefficient dynamics agree with computing shear coordinates with laminations.

Thus, we can use laminations to build an arbitrary cluster algebra of geometric type (with $(m+n) \times n$ exchange matrix) as long as $n \times n$ top corresponds to a cluster algebra from a surface.

Thm (w. Thurston) For a fixed triangulation T without self-folded triangles, the map

$$L \rightarrow (b_E(T, L))_{E \in T} \text{ is a}$$

bijection between integral unbounded measured laminations and \mathbb{Z}^n .

(13) Consequently, if T has no self-folded triangles, any desired coeff. pattern can be achieved by m laminations. Each L_i corresponds to a row of coefficients.

Thm (w. Thurston, Fock, Goncharov)

If T, T' are triangulations without self-folded triangles are related by flipping edge E_k , then $(m+n) \times n$ exch. matrices

$\tilde{B}(T, L)$ & $\tilde{B}(T', L)$ are related by M_k .

Fomin-Thurston defined shear coordinates for triangulations with self-folded triangles or tagged arcs too.

Enlarges bijection and Theorem.

Next time: Lambda Lengths,
a Teichmüller interpretation of cluster variables.