## Super Cluster Algebras from Surfaces

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$$
\begin{aligned}
& \text { arXiv:2102.09143 } \\
& \text { arXiv:2110.06497 } \\
& \operatorname{arXiv:2208.13664}
\end{aligned}
$$



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## Background

## Super Cluster Algebras from Surfaces

Let $F$ be a bordered surface with marked points on its boundary. Loosely speaking, a cluster algebra from $F$ has...
(1) clusters $\Longleftrightarrow$ ideal triangulations of $F$
(2) cluster variables $\Longleftrightarrow$ "lengths" of diagonals
(3) mutations $\Longleftrightarrow$ Ptolemy relations


$$
e f=a b+c d
$$

"Super" means super-commutative, i.e.

$$
A=A_{0} \oplus A_{1} \quad \text { with relations } \quad x y=(-1)^{\bar{x} \bar{y}} y x
$$

where $\bar{x}=i$ if $x \in A_{i}$.
More specifically, for $a, b \in A_{0}$ and $\theta, \sigma \in A_{1}$ we have

$$
a b=b a \quad a \theta=\theta a \quad \theta \sigma=-\sigma \theta
$$

## Question

Define a super-commutative analogue of cluster algebras?

- We will often work with a superalgebra $A=A_{0} \oplus A_{1}$.
- Elements in $A_{0}$ are commutative, which are called even or bosonic, and will be denoted by Latin letters $x, y, z \cdots$.
- Elements in $A_{1}$ are anti-commutative, which are called odd or fermionic, and will be denoted by Greek letters $\theta, \sigma, \alpha, \beta, \cdots$.
- An element in a superalgebra has a body and a soul...

$$
\underbrace{1+x_{1} x_{2}+x_{3}}_{\text {body }}+\underbrace{x_{1} \theta_{1} \theta_{2}+\theta_{1}+\left(x_{1}-x_{2}\right) \theta_{2}}_{\text {soul }}
$$

- An important fact is that odd variables square to zero: $\theta^{2}=0$.
(1) Motivation
(2) Decorated Super Teichmüller Theory
(3) First Formula: Super $T$-paths

4. Second Formula: Double Dimers
(5) $\operatorname{OSp}(1 \mid 2)$-Matrix Formula
(6) Super Fibonacci Numbers

## Outline

(1) Motivation
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## A Brief History of Cluster Superalgebras

(1) Ovsienko proposed an approach for cluster superalgebras [Ovs15] motivated by the study of superfriezes [MGOT15].
(2) This approached was later expanded to the definition of cluster algebras with Grassmann variables by Ovsienko-Shapiro [OS18].
(3 [LMRS17] gives a different approach, based on superfrieze patterns and $\operatorname{Gr}(2|0,4| 1)$.
44 [SV19] computed super Plüker relation for super Grassmannians and discussed certain cluster structures there-in. A more detailed discussion for the case $\operatorname{Gr}(2,0 \mid n, 1)$ was given in [She22] very recently.
© In [MOZ21, MOZ22a, MOZ22b], Musiker-Ovenhouse-Z. studied the cluster structure of Penner-Zeitlin's decorated super-Teichmüller spaces.

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## Decorated Teichmüller Theory

The Teichmüller space of a surface $F=F_{g}^{s}$ is

$$
T(F)=\operatorname{Hom}\left(\pi_{1}(F), \operatorname{PSL}(2, \mathbb{R})\right) / \operatorname{PSL}(2, \mathbb{R})
$$

And the decorated Teichmüller space is the trivial $\mathbb{R}_{>0}^{s}$-bundle over $T(F)$, denoted $\tilde{T}(F)$. See [Pen87].
Roughly speaking, there is a $\lambda$-length associated to every pair of ideal points, satisfying the Ptolemy relation:

where $e f=a c+b d$.

## Decorated Super-Teichmüller Spaces

- By replacing $\operatorname{PSL}(2, \mathbb{R})$ with $\operatorname{OSp}(1 \mid 2)$, the super-Teichmüller space of a surface $F$ is

$$
S T(F)=\operatorname{Hom}\left(\pi_{1}(F), \operatorname{OSp}(1 \mid 2)\right) / \operatorname{OSp}(1 \mid 2)
$$

- In the decorated space, we have, similar to the classical case, a super $\lambda$-length for every pair of ideal points; and
- new coordinates called $\mu$-invariants for every triple of ideal points (i.e. triangles).
- In addition, the super Teichmüller space consists of connected components indexed by spin structures, which are equivalence classes of orientations on the triangulations.



## Super Ptolemy Relation

The Ptolemy transformation on super $\lambda$-length coordinates is given as follows.


$$
\begin{aligned}
e f & =a c+b d+\sqrt{a b c d} \sigma \theta \\
\sigma^{\prime} & =\frac{\sigma \sqrt{b d}-\theta \sqrt{a c}}{\sqrt{a c+b d}} \text { and } \theta^{\prime}=\frac{\theta \sqrt{b d}+\sigma \sqrt{a c}}{\sqrt{a c+b d}} \\
\sigma \theta & =\sigma^{\prime} \theta^{\prime}
\end{aligned}
$$

## Super Ptolemy Relation

Super-flip reverse the orientation of the edge $b$.


## Remark

- Super Ptolemy moves are not involution: $\mu_{i}^{8}=I$.
- The body of a super $\lambda$-length are exactly the (ordinary) $\lambda$-length in the bosonic $T(F)$.

If we flip a diagonal twice:


The orientations of the triangle $\theta$ are reversed and $\theta$ is changed to $-\theta$, which corresponds to the equivalence relation mentioned before. In other words, super Ptolemy relations are involutions only up to equivalence.

## Super Ptolemy Relation - Example



Start with a Pentagon with given orientation, and we will calculate the super $\lambda$-length of the longest diagonal by flipping $x_{1}$ then $x_{2}$.

We first flip the edge $x_{1}$.

## Super Ptolemy Relation - Example

After flipping $x_{1}$ to $x_{3}$, we get:

$$
\begin{aligned}
& x_{3}=\frac{a d+e x_{2}}{x_{1}}+\frac{\sqrt{a d e x_{2}}}{x_{1}} \theta_{1} \theta_{2} \\
& \theta_{4}=\frac{\sqrt{a d} \theta_{1}-\sqrt{e x_{2}} \theta_{2}}{\sqrt{x_{1} x_{3}}} \\
& \theta_{5}=\frac{\sqrt{a d} \theta_{2}+\sqrt{e x_{2}} \theta_{1}}{\sqrt{x_{1} x_{3}}}
\end{aligned}
$$

Here the red color indicates that the orientation has been reversed.

Next we flip $x_{2}$.

## Super Ptolemy Relation - Example

After flipping $x_{2}$ to $x_{4}$, we have:

$$
x_{4}=\frac{a c+b x_{3}}{x_{2}}+\frac{\sqrt{a c b x_{3}}}{x_{2}} \theta_{5} \theta_{3}
$$

$$
=\frac{a c x_{1}+a b d+b e x_{2}}{x_{1} x_{2}}+\frac{b \sqrt{a d e x_{2}}}{x_{1} x_{2}} \theta_{1} \theta_{2}+
$$

$$
\frac{\sqrt{a c b\left(\frac{a d+e x_{2}}{x_{1}}+\frac{\sqrt{a d e x_{2}}}{x_{1}} \theta_{1} \theta_{2}\right)}}{x_{2}}\left(\frac{\sqrt{a d} \theta_{2}+\sqrt{e x_{2}} \theta_{1}}{\sqrt{x_{1} x_{3}}}\right) \theta_{3}
$$

$$
=\frac{a c x_{1}}{x_{1} x_{2}}+\frac{a b d}{x_{1} x_{2}}+\frac{b e x_{2}}{x_{1} x_{2}}+\frac{b \sqrt{a d e}}{x_{1} \sqrt{x_{2}}} \theta_{1} \theta_{2}+
$$

$$
\frac{a \sqrt{b c d}}{\sqrt{x_{1} x_{2}}} \theta_{2} \theta_{3}+\frac{\sqrt{a b c d}}{\sqrt{x_{1} x_{2}}} \theta_{1} \theta_{3}
$$

## Question

In a cluster algebra $A$, any cluster variable can be expressed as a positive Laurent polynomial in the initial cluster, i.e.

$$
A \subset \mathbb{R}\left[x_{1}^{ \pm 1}, \cdots, x_{n}^{ \pm 1}\right]
$$

## Questions

- Does the super $\lambda$-length satisfy some Laurent phenomenon?
- Is there a "positivity" for terms with anti-commuting variables?


## Answers (Spoiler Alert)

- Super $\lambda$-lengths live in $\mathbb{R}\left[x_{1}^{ \pm \frac{1}{2}}, \cdots, \left.x_{1}^{ \pm \frac{1}{2}} \right\rvert\, \theta_{1}, \cdots, \theta_{n+1}\right]$.
- There exists an ordering on the odd variables, called positive ordering, such that if we multiply $\theta^{\prime}$ s in the positive ordering then the coefficients are positive.


## Modified $\mu$-invariants

Now we introduce some new notations to simplify the calculations.
For a triangle


Define the $h$-lengths

$$
h_{j k}^{i}=\frac{\lambda_{j k}}{\lambda_{i j} \lambda_{i k}}, h_{i k}^{j}=\frac{\lambda_{i k}}{\lambda_{i j} \lambda_{j k}}, h_{i j}^{k}=\frac{\lambda_{i j}}{\lambda_{i k} \lambda_{k j}}
$$

and

$$
\begin{aligned}
& \Delta_{j k}^{i}:=\sqrt{\frac{\lambda_{j k}}{\lambda_{i j} \lambda_{i k}}} \theta=\sqrt{h_{j k}^{i}} \theta, \Delta_{i k}^{j}:=\sqrt{h_{i k}^{j}} \theta, \Delta_{i j}^{k}:=\sqrt{h_{i j}^{k}} \theta, \\
& \nabla_{j k}^{i}:=\sqrt{\frac{\lambda_{i j} \lambda_{i k}}{\lambda_{j k}}} \theta=\sqrt{\frac{1}{h_{j k}^{i}}} \theta, \nabla_{i k}^{j}:=\sqrt{\frac{1}{h_{i k}^{j}}} \theta, \nabla_{i j}^{k}:=\sqrt{\frac{1}{h_{i j}^{k}}} \theta .
\end{aligned}
$$

## Super Ptolemy Relations Revisited




From now on, only consider triangulations with a longest diagonal, and decompose into fans whose centers are labelled $c_{1}, c_{2}, \cdots$.
Define a default orientation as follows

- Edges inside each fan segments are directed away from the center.
- Others are oriented as

$$
c_{1} \rightarrow c_{2} \rightarrow \cdots \rightarrow c_{n}
$$

Define a positive ordering on $\mu$-invariants.

- Going from bottom to top, append the odd variable to the left (resp. right) if the arrow is pointing left (resp. right).
$\alpha_{1}>\alpha_{2}>\alpha_{3}>\gamma_{1}>\gamma_{2}>\gamma_{3}>\delta_{2}>\delta_{1}>\beta_{2}>\beta_{1}$


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## Review of Schiffler's (ordinary) T-paths

A $T$-path from $i$ to $j$ is a path on the triangulation $T$ starting at vertex $i$, ending at $j$, such that
(T1) the path does not use any edge twice
(T2) the path has an odd number of edges
(T3) the even-numbered edges cross the diagonal $(i, j)$
(T4) The path is getting closer from $i$ to $j$.
Assign a $T$-path a weight to $\mathrm{wt}(t)=\frac{\Pi \text { odd edges }}{\Pi \text { even edges }}$, then the cluster variable ( $\lambda$-length) $\lambda_{i j}$ is the weighted sum of all $T$-paths from $i$ to $j$.

$$
\frac{x_{23} x_{15}}{x_{13}}
$$



$$
\frac{x_{12} x_{34} x_{15}}{x_{13} x_{14}}
$$


$\lambda_{25}=\sum_{t \in T_{25}} \mathrm{wt}(t)=\frac{x_{23} x_{15}}{x_{13}}+\frac{x_{12} x_{34} x_{15}}{x_{13} x_{14}}+\frac{x_{12} x_{45}}{x_{14}}$


Super T-paths are paths on the auxiliary graph, where all the usual T-paths moves are allowed.

The additional moves are

- Enter or leave the internal (only) at odd steps, with wt $\left({\underset{j}{ }}_{\dot{\delta}}^{k}\right)=\Delta_{j k}^{i}$.
- Can teleport from an internal vertex to another, with weight 1 .


## Theorem (Musiker-Ovenhouse-Z. 21)

For a default orientation, super $\lambda$-lengths are (positive) weighted sums of super $T$-paths, where all products of odd variables are written in the positive ordering.

## Super T-paths: Examples



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## Formula for $\lambda$-lengths: Example



$$
\theta_{1}>\theta_{2}>\theta_{3}
$$



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## Snake Graphs

Ordinary cluster variables can also be seen as perfect matchings (dimer covers) of snake graphs.


## Dimer Covers on Snake Graphs

A dimer cover (a.k.a perfect matching) $M$ of a graph $G$ is a collection of edges such that every vertex in $G$ is incident to exactly one edge in $M$.

The weight of a dimer cover is the product of the edge weights.


$$
\text { weight }=b f d x y z
$$

## Theorem (Musiker-Schiffler, Musiker-Schiffler-Williams)

The $\lambda$-length is the given by

$$
\lambda(\gamma)=\frac{1}{\operatorname{cross}(\gamma)} \sum_{\substack{M \text { dimer cover } \\ \text { of the snake graph }}} \mathrm{wt}(M)
$$

## Double Dimer Covers

Surprisingly, the super $\lambda$-lengths naturally arise as double dimer covers of the same snake graph, which are unions of two dimer covers and contains single edges and doubled edges.
The weight of a double dimer cover is the product of the square root of it edges, multiplied by the odd variables on the first and last triangle of cycles.


$$
\text { weight }=x y z \sqrt{a b c d e f} \theta_{1} \theta_{3}
$$

Theorem (Musiker-Ovenhouse-Z. 22a)
The $\lambda$-length is the given by

$$
\lambda(\gamma)=\frac{1}{\operatorname{cross}(\gamma)} \sum_{\substack{M \text { dimer cover } \\ \text { of the sakake oranh }}} w t(M)
$$

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## $\operatorname{OSp}(1 \mid 2)$

The orthosymplectic supergroup $\operatorname{OSp}(1 \mid 2)$ contains the set of $2|1 \times 2| 1$ matrices

$$
M=\left(\begin{array}{ll|l}
a & b & \gamma \\
c & d & \delta \\
\hline \alpha & \beta & e
\end{array}\right)
$$

such that

$$
\begin{array}{lll}
e=1+\alpha \beta & e^{-1}=a d-b c & \alpha=c \gamma-a \delta \\
\beta=d \gamma-b \delta & \gamma=a \beta-b \alpha & \delta=c \beta-d \alpha
\end{array}
$$

Note that it contains a $\mathrm{SL}_{2}$ subgroup

$$
\left(\begin{array}{ll|l}
a & b & 0 \\
c & d & 0 \\
\hline 0 & 0 & 1
\end{array}\right)
$$

## Special elements of $\operatorname{OSp}(1 \mid 2)$

Let $x, h$ be even and $\theta$ odd, we define

$$
\begin{gathered}
E(x)=\left(\begin{array}{cc|c}
0 & -x & 0 \\
1 / x & 0 & 0 \\
\hline 0 & 0 & 1
\end{array}\right) \quad A(h \mid \theta)=\left(\begin{array}{cc|c}
1 & 0 & 0 \\
h & 1 & -\sqrt{h} \theta \\
\hline \sqrt{h} \theta & 0 & 1
\end{array}\right) \\
\rho=\left(\begin{array}{cc|c}
-1 & 0 & 0 \\
0 & -1 & 0 \\
\hline 0 & 0 & 1
\end{array}\right)
\end{gathered}
$$

Their inverses are given by $\rho^{-1}=\rho, E(x)^{-1}=\rho E(x)=E(-x)$ and

$$
A(h \mid \theta)^{-1}=\left(\begin{array}{cc|c}
1 & 0 & 0 \\
-h & 1 & \sqrt{h} \theta \\
\hline-\sqrt{h} \theta & 0 & 1
\end{array}\right)
$$

Note that $\rho A(h \mid \theta) \rho=A(h \mid-\theta)$. This corresponds to the equivalence relation of orientations in a spin structure.
We will also abbreviate

$$
E_{i j}:=E\left(\lambda_{i j}\right)=\left(\begin{array}{cc|c}
0 & -\lambda_{i j} & 0 \\
\lambda_{i j}^{-1} & 0 & 0 \\
\hline 0 & 0 & 1
\end{array}\right) \quad A_{j k}^{i}:=A\left(h_{j k}^{i} \mid \overleftarrow{i j k}\right)=\left(\begin{array}{cc|c}
1 & 0 & 0 \\
h_{j k}^{i} & 1 & -\Delta_{j k}^{i} \\
\hline \Delta_{j k}^{i} & 0 & 1
\end{array}\right)
$$

## A graph on $T$

From a triangulation $T$ of a marked surface, we associate a graph $\Gamma_{T}$ by putting 6 vertices inside each triangle, and connect them in the following way


Figure: The graph $\Gamma_{T}$, with $T$ in dashed lines.

For a graph embedded on a surface, a graph connection is an assignment of a matrix to each oriented edge, such that the opposite oriented edge are assigned to its inverse.
For a path in the graph, the holonomy is the corresponding product of matrices along the path.
If the path is a loop, then the holonomy is also called monodromy.
A connection is called flat if the monodromy around each contractible face is the identity matrix.

## A Flat $\operatorname{OSp}(1 \mid 2)$-connection on $\Gamma_{T}$.

For each oriented edge of $\Gamma_{T}$, associate an element of $\operatorname{OSp}(1 \mid 2$ as follows.

| Type (i) |  | $A_{j k}^{i}{ }^{-1}$ |  | $A_{j k}^{i}$ |
| :---: | :---: | :---: | :---: | :---: |
| Type (ii) |  | $E_{i j}^{-1}$ |  | $E_{i j}$ |
| Type (iii) |  | $\rho$ |  | id |

This defines a flat $\operatorname{OSp}(1 \mid 2)$-connection on $\Gamma_{T}$.

## Matrix Formula for super $\lambda$-length



The holonomy matrix from a point near $i$ to a point near $k$ is given by

$$
H_{i k}=\left(\begin{array}{cc|c}
-\frac{\lambda_{j k}}{\lambda_{i j}} & \pm \lambda_{i k} & \nabla_{i j}^{k} \\
\pm \frac{\lambda_{j l}}{\lambda_{i j} \lambda_{k l}} & \pm \frac{\lambda_{k l}}{\lambda_{k l}} & \pm \frac{1}{\lambda_{k l}} \nabla_{i j}^{l} \\
\hline \frac{1}{\lambda_{i j}} \nabla_{k l}^{j} & \pm \nabla_{k l}^{i} & 1+\star
\end{array}\right)
$$

In particular, the $(2,2)$-entry is the super $\lambda$-length up to sign.

The proof uses induction in two different ways, by left-multiplication and right-multiplication.

By induction via left-multiplication, we prove the first two columns, which corresponds to flipping the diagonals from bottom to top.

By induction via right-multiplication, we prove the first two rows, which corresponds to flipping the diagonals from top to bottom.

## Connection to Double Dimers

The following matrix, whose entries are weighted sum of certain double dimer covers, satisfies the $\operatorname{OSp}(1 \mid 2)$ relations.


This is an analogue of 'Kuo's condensation'.
(1) The $\mathrm{SL}_{2}$ part of our matrix formula is the same as the one given by Musiker-Williams up to signs. In particular, the usage of $E$ and $E^{-1}$ are swapped.
(2) A similar construction for sheer coordinates of super Teichmüller spaces was given by F. Bouschbacher in his thesis. In cluster algebra language, shear coordinates are $\mathcal{X}$-type cluster variables, while $\lambda$-lengths are $\mathcal{A}$-type cluster variables.
(3) The constructions given for $\Gamma_{T}$ and the connection make sense for any triangulated surface. For a surface with non-trivial topology, the monodromy of this connection coincide with the representation $\pi_{1}(S) \rightarrow \operatorname{OSp}(1 \mid 2)$ described in Section 6 of Penner-Zeitlin.

## Super Fibonacci Numbers

Consider an annulus with one marked point on each boundary component, and the oriented triangulation, where all $\lambda$-lengths are equal to 1 .
Let $z_{n}$ be $\lambda$-length of the arc connecting the two marked points which winds around the annulus $n-2$ times. This is the analogue of even indexed Fibonacci number.


In our previous paper, we showed that

$$
z_{n}=(3+2 \sigma \theta) z_{n-1}-z_{n-2}-\sigma \theta,
$$

## Super Fibonacci Numbers Continued

Let $z_{n}=x_{2 n-5}+y_{2 n-5} \sigma \theta$ and define $w_{n}=x_{2 n-4}+y_{2 n-4} \sigma \theta$, they satisfy the following recurrence.
(a) $z_{n}=z_{n-1}+(1+\sigma \theta) w_{n-1}$
(b) $w_{n}=w_{n-1}+(1+\sigma \theta) z_{n}-\sigma \theta$

By means of our matrix formula, we now give an interpretation for the $w_{n}{ }^{\prime}$ s.
$H\left(z_{n}\right)=\left(\begin{array}{cc|c}-w_{n-1} & z_{n} & \left(z_{n}-1\right) \sigma+w_{n-1} \theta \\ -z_{n-1} & w_{n-1} & \left(z_{n-1}-1\right) \theta+w_{n-1} \sigma \\ \hline\left(z_{n-1}-1\right) \sigma-w_{n-1} \theta & \left(z_{n}-1\right) \theta-w_{n-1} \sigma & 1-\left(\ell_{2 n-4}-2\right) \sigma \theta\end{array}\right)$
where $\ell_{n}$ is the Lucas number.

## Thank you for listening!

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